

Toward single e-bunch shape diagnostics using THz coherent radiation

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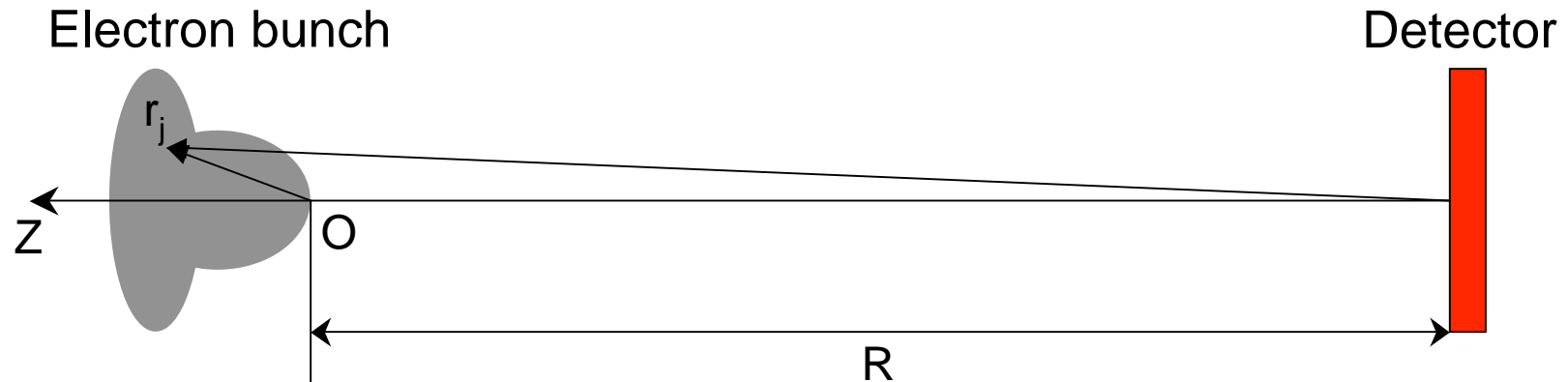
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Outline

- Coherent radiation and the electron bunch shape
- Holographic spectroscopy in the THz
- Bringing measurement into the visible with EO detection
- Timeline

Coherent emission of a distant bunch



Total intensity of the bunch radiation

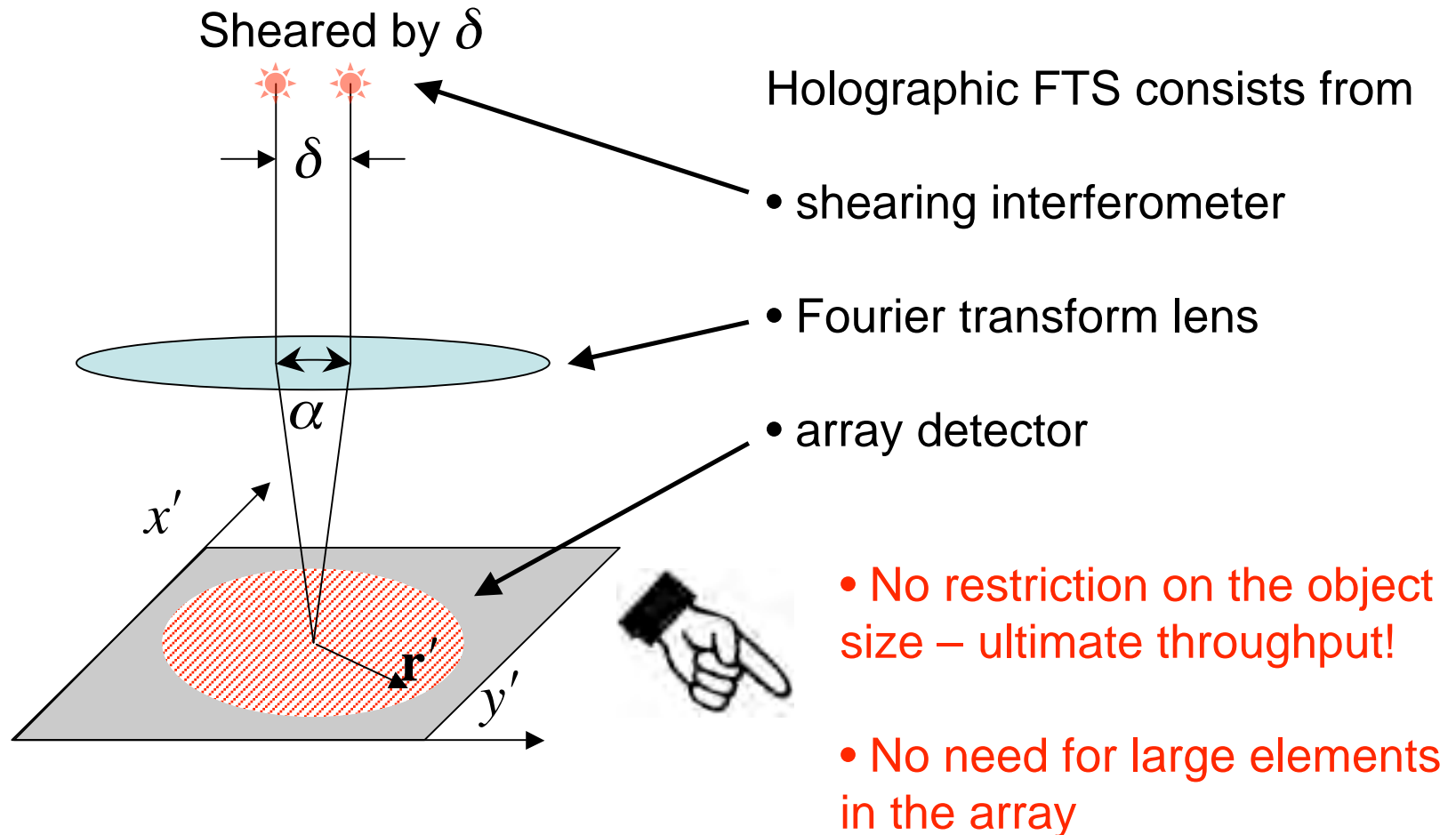
$$I_{tot}(\omega) = I(\omega)[N + N(N-1)F(\omega)]$$

Electron bunch form-factor

$$F(\omega) = \left| \int_0^\infty dz S(z) e^{i(\omega/c)z} \right|^2$$

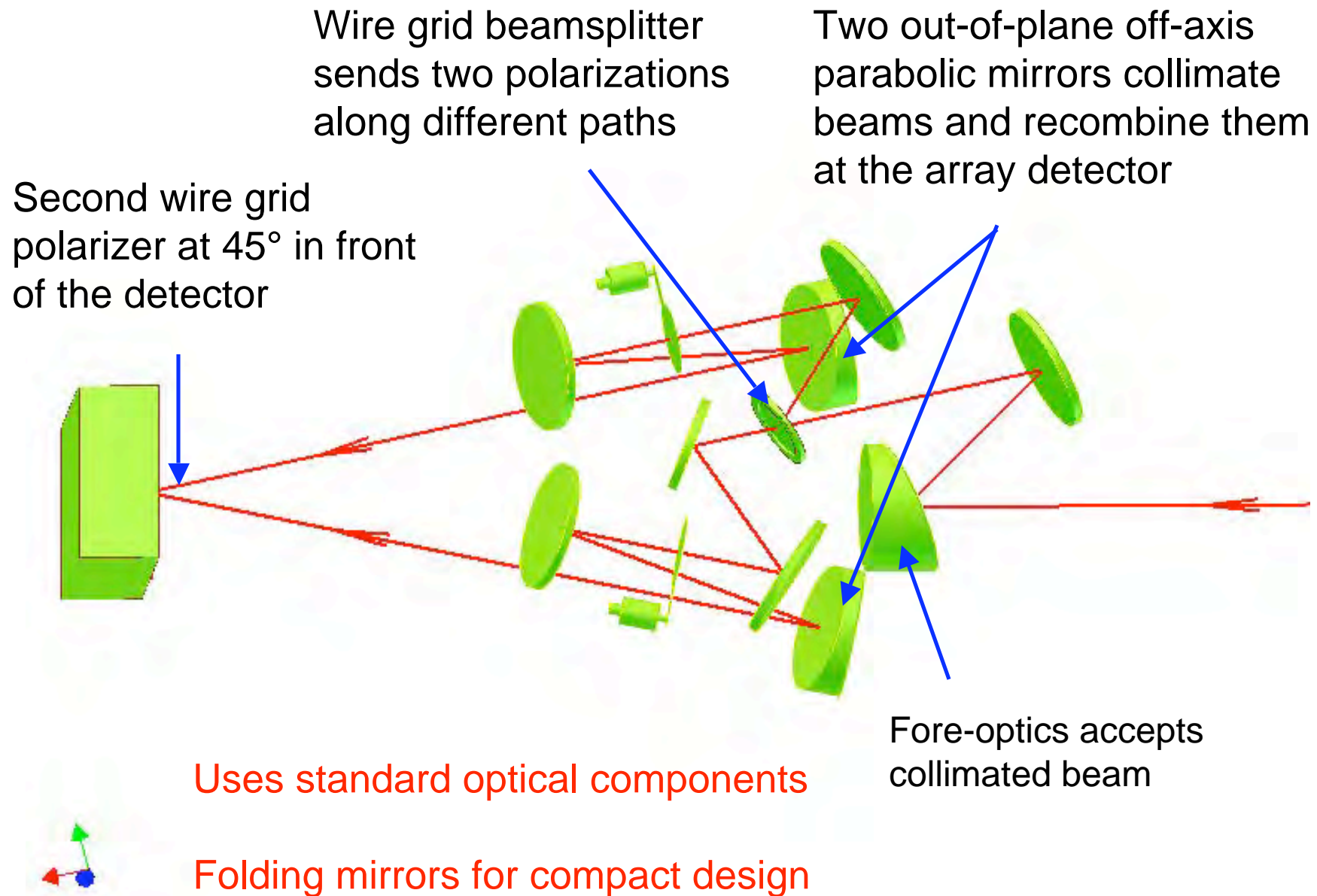
Coherent radiation spectrum contains information
about the bunch shape

Properties of ideal HFTS

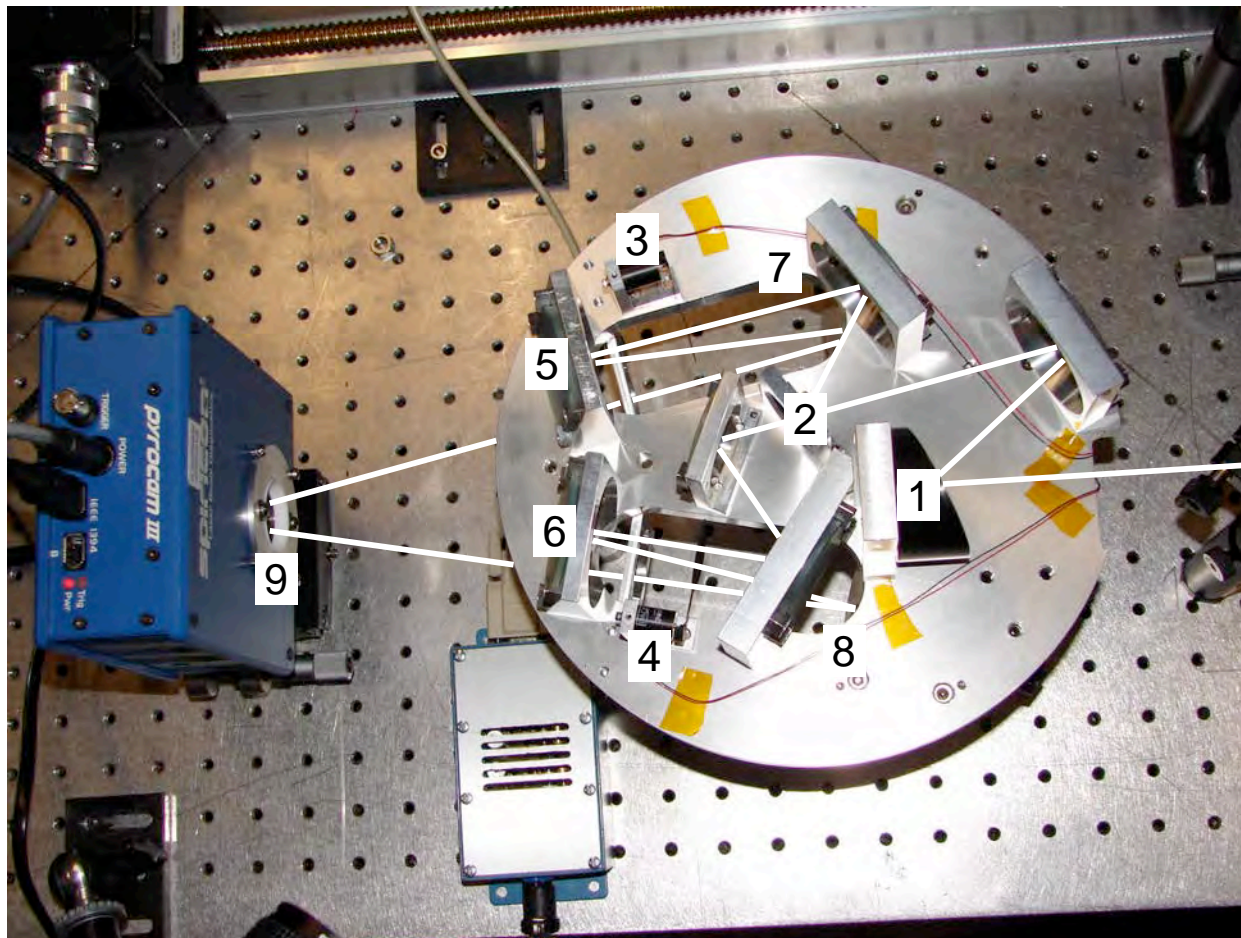


$$I(\mathbf{r}') = 1/2 \int_0^\infty d\omega I_0(\mathbf{r}', \omega) \left(1 + \cos \frac{\omega}{c} \frac{\delta}{f} y' \right) B(\omega)$$

Optical layout of the THz interferometer



Experimental setup at J-Lab



1 - fore-optics

2 - beamsplitter

3,4 - shutters

5,6 - tilted folding
mirrors

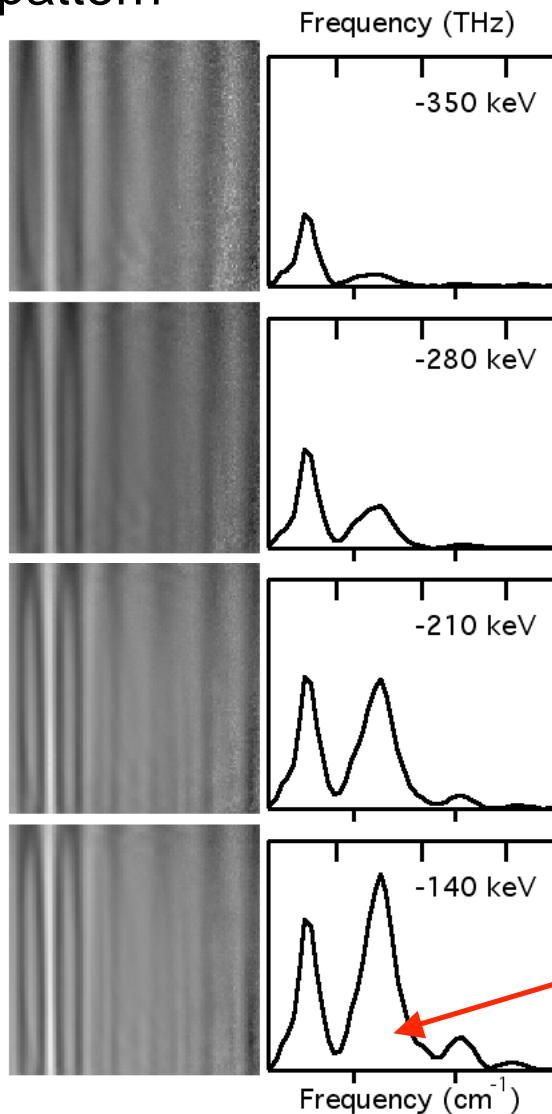
7,8 - off-axis parabolic
mirrors

9 - polarizer

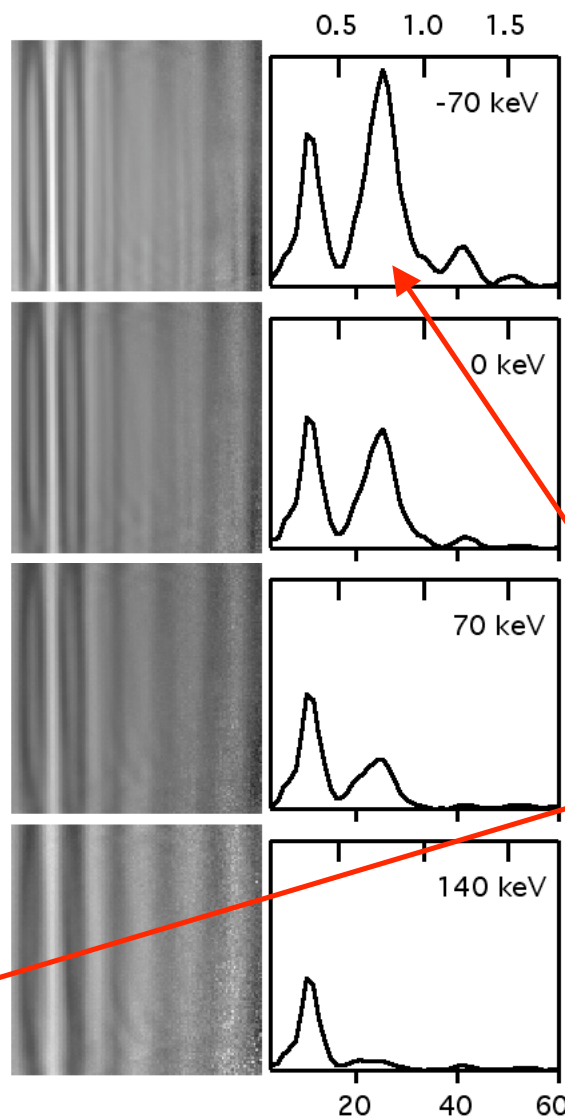
Experiments conducted in air since only a few water lines occur below 40 cm^{-1}

Coherent SR measurements at JLab

Interference
pattern



Spectrum



ERL parameters:

current - 1 mA

charge - 100 pC

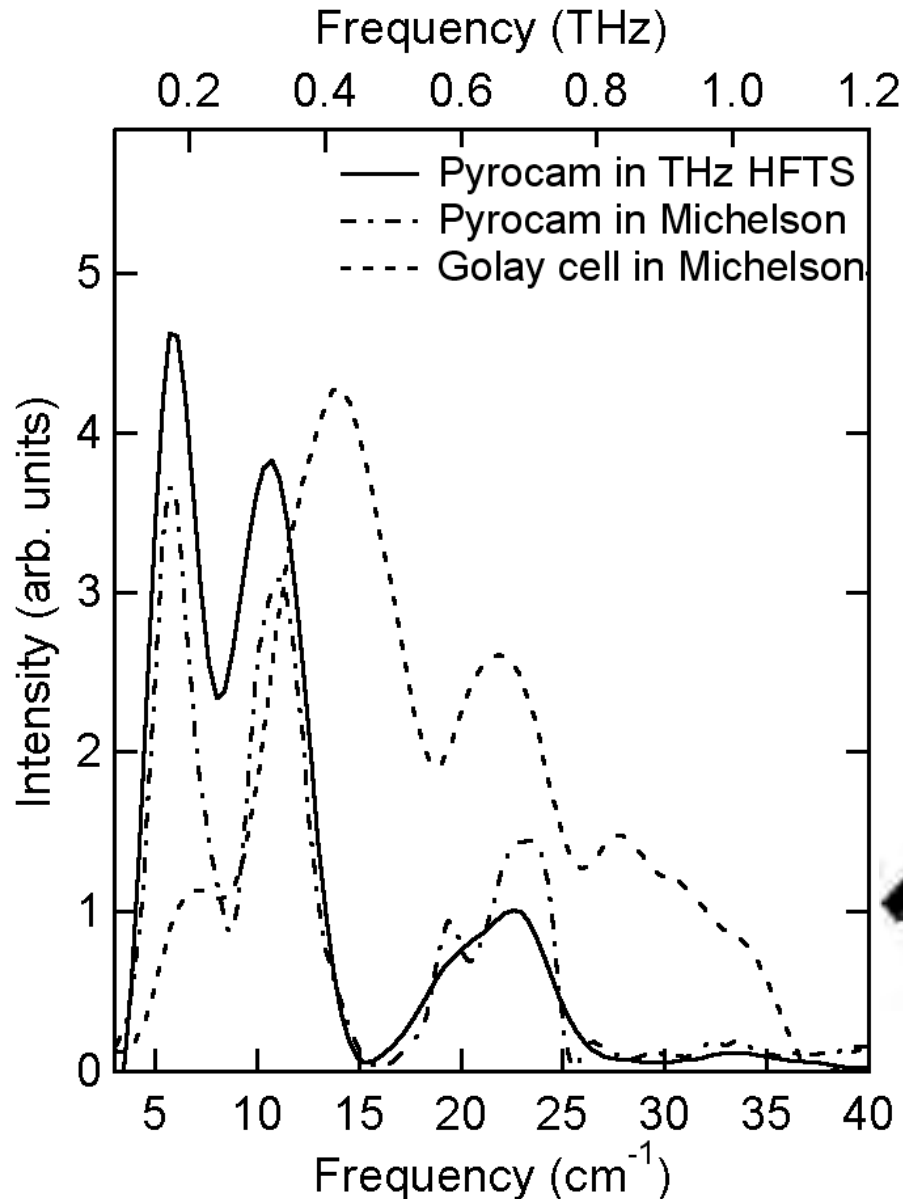
frequency - 9.36 MHz

energy - 114.65 MeV

Detuning of the bunch
energy by changing the
gradient field in the RF
cavities varied the electron
bunch length

The stronger and
broader spectra are
observed for shorter
bunches

Non-uniform sensitivity of Pyrocam array



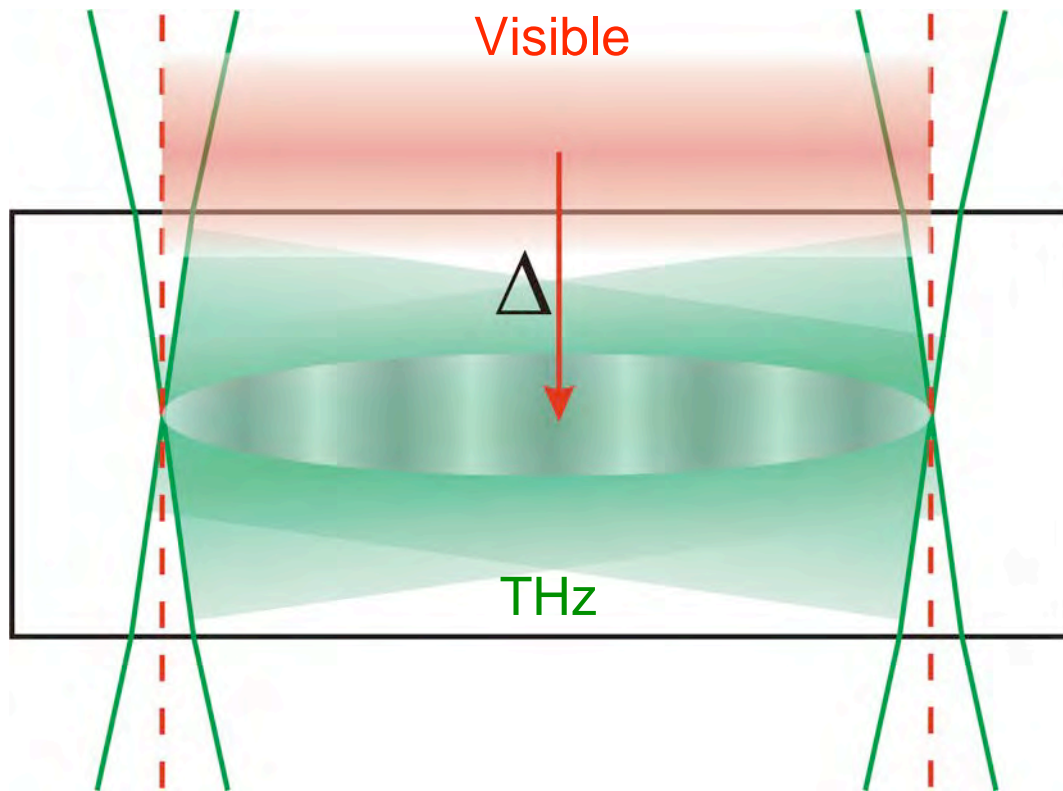
Three measurements were performed at JLab

1. THz HFTS with Pyrocam array
2. Scanning Michelson FTS with the Pyrocam
3. Scanning Michelson FTS with the Golay cell



Dips in the spectra are due to the Pyrocam array non-uniform sensitivity

Electro-optic detection of the THz interference



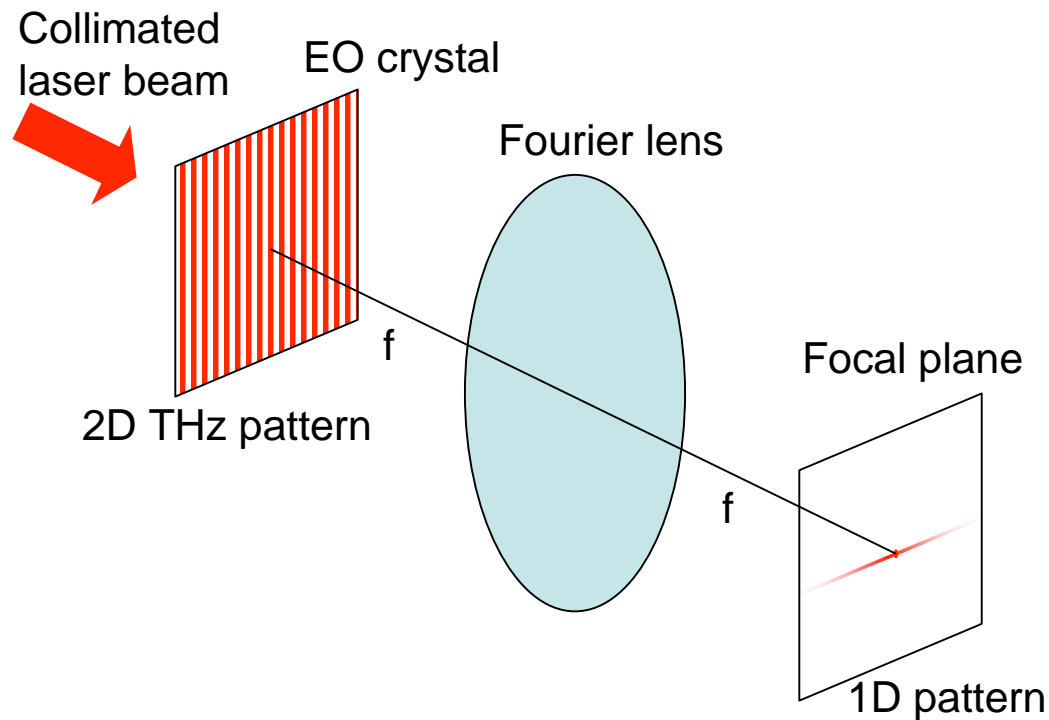
Delay Δ between THz and visible pulses is variable

Existing THz arrays insensitive and not optimized for spectral measurements

The way forward: use electro-optic detection

- EO effect is fast
- No pixel size restriction
- Flexibility and expandability

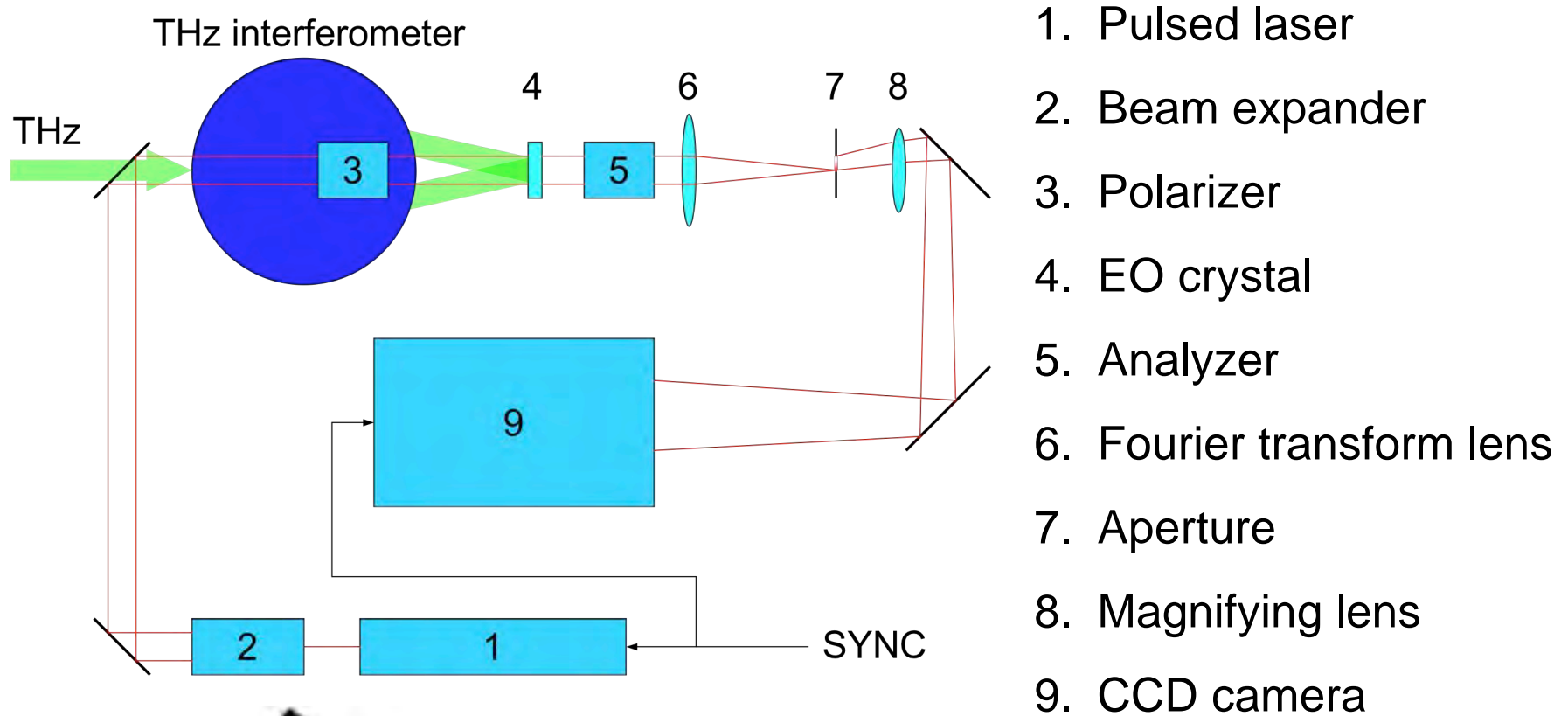
Optical FT of the THz interference pattern



Optical FT advantages

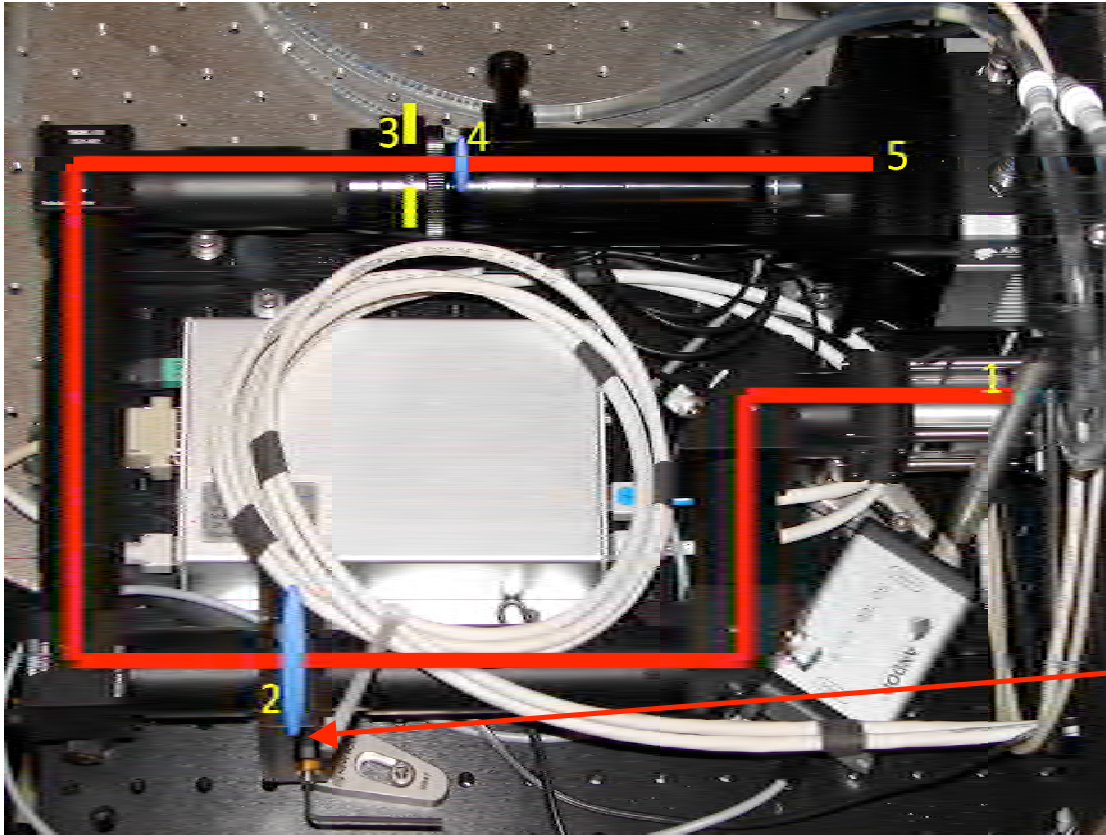
- Signal concentrated in a 1D pattern increasing the signal-to-noise ratio
- No signal in the absence of THz radiation - no need to use ultrafast laser pulses
- Rejection of the DC component - decreased background

Schematic of the THz HFTS with optical FT



A self contained instrument without
external synchronized lasers

EO detection module for optical FT



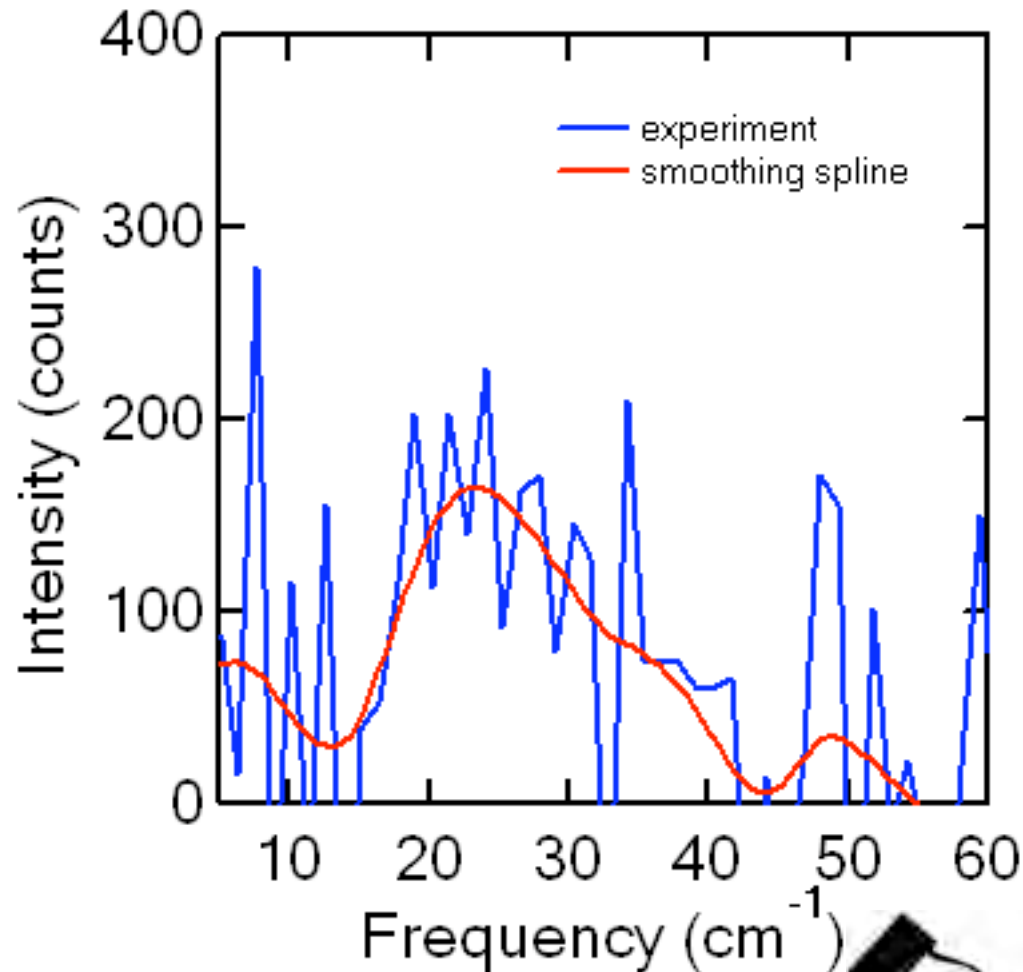
1. Analyzer
2. FT lens
3. Aperture
4. Magnifying lens
5. CCD camera

The FT lens is mounted on the XY stage in order to move the DC component outside of the aperture (3)



Complete instrument including the diode laser, THz interferometer, and the EO detection module fits on the 12" X 36" optical board

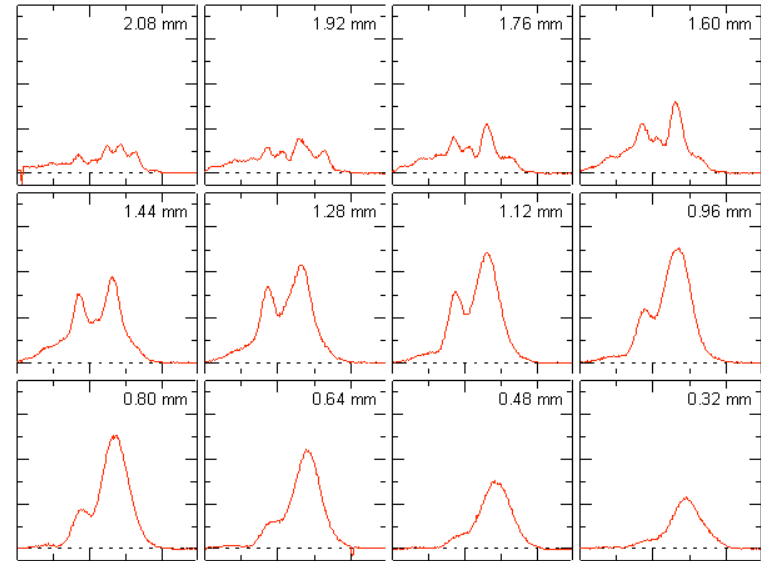
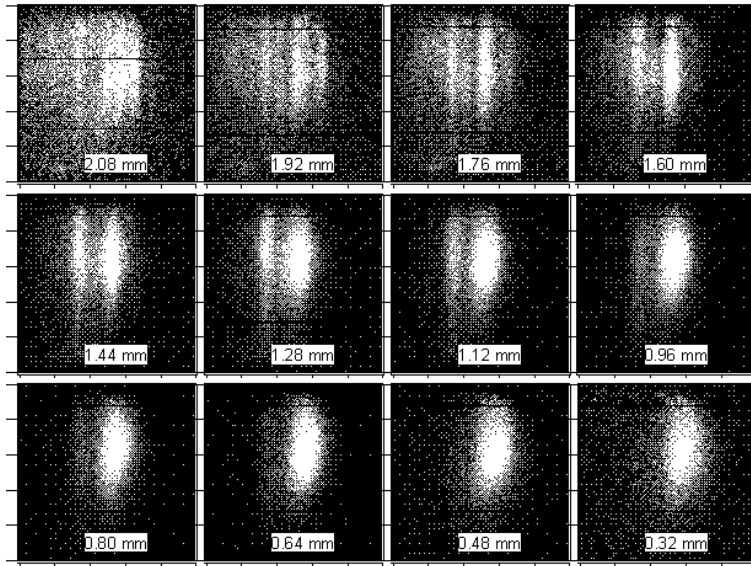
Experimental results for HFTS with optical FT at DESY



- Operation of the THz interferometer in the imaging mode
- Large background due to the long diode laser pulse and the EO crystal inhomogeneity
- Weak signal for the crossed polarizers geometry
- Accumulation of multiple shots

For single shot capability needed laser with < 100 ps pulses, > 1 W peak power

THz interference readout with incoherent radiation at JLab

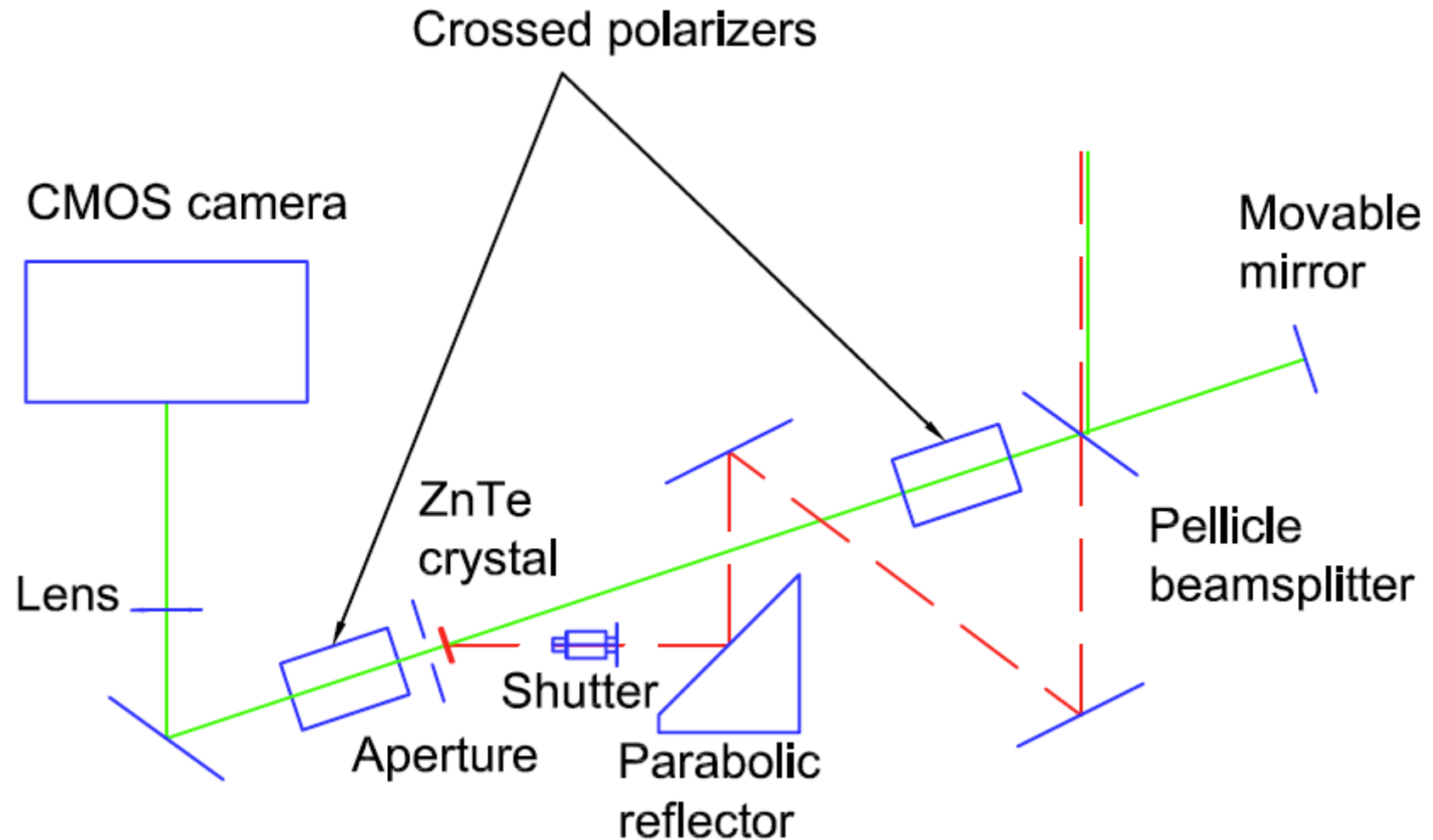


Frames (left) and corresponding profiles (right) reflect change in the THz interference pattern with the varying delay Δ between the incoherent and coherent pulses



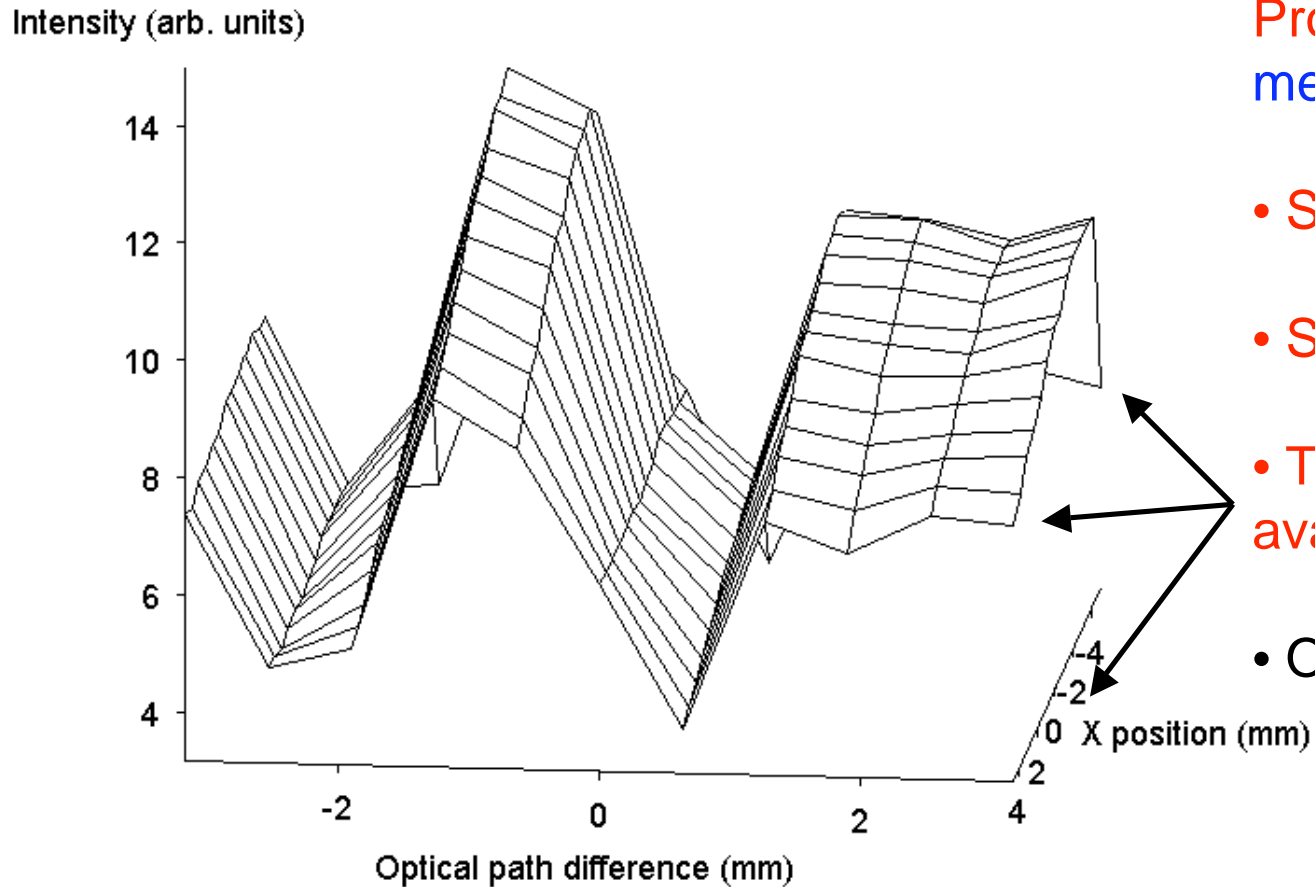
- Self-contained, no external laser
- Single shot capable
- Asymmetry capable

Cross-correlation of incoherent and coherent radiation



Simple, compact and self-contained system

Cross-correlation results at JLab



Pros and cons of the method

- Simple, self-contained
 - Single shot capable
 - Transverse dimension available as well
 - Only symmetric shape
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- Preliminary results confirm the idea
 - Effects of crystal dispersion on the EO response function should be investigated



Timeline

Past
multi shot

Present
multi shot

Future
single shot

