

Toward Extreme-Scale: High-Order Algorithms for Electromagnetic and Fluid Modeling

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Outline

- Overview for Application Problems
- Motivation for using High-Order Methods
- Production Codes: NekCEM and NekLBM
- **Efficiency and Scalability on >100K cores**
- Applications: Results on Nanoscience, Accelerator, and Fluid Modeling
- **Summary for Future Plans**

Applications

Introduction



Motivation for High-Order Methods

- Representing geometry accurately: material discontinuities, conforming meshes
- Long time integration requires minimal numerical dispersion
- High-order methods deliver engineering accuracy with fewer points per wavelength.
 (min ppw = 2 by Nyquist sampling thm)



- Efficient implementation
- Scalable parallel performance
- Flexible geometry representation



1D Maxwell, n=120, varying wavenumber

NekCEM/NekLBM

- Spectral-Element DG Time-Domain Solvers: using core structure of *Nek5000*
 - Open source code: https://svn.mcs.anl.gov/repos/NEKCEM
 - Open source code: https://svn.mcs.anl.gov/repos/NEKLBM
 - Argonne developed (initiated) Maxwell and Boltzmann solvers in Fortran and C.
 - Currently scalable up to > 100K cores (more than 1.1 B grids points)
 - Paralle I/O: output in a single or multiple files: scalable up to 65K cores



NekCEM Runtime Tasks

Parallel Scalability



Excellent scaling results from:

- DG: only six messages per element (not 26)
- Vector communication (1 msg for 6 fields)

Parallel I/O Scalability

Software Development

ICS 2011, Jng Fu, RPI



rbIO: reduced-blocking I/O coIO: collective I/O

Parallel Scalability, LBM Code



Excellent scaling results from:

- DG: only six messages per element (not 26)

- Vector communication (1 msg for 19 fields)

Weak Formulation

□ Maxwell Equation (source-free in free space) and weak form:

$$\Box \frac{\partial q}{\partial t} + \nabla \bullet F(q) = 0 \xrightarrow{\text{weak form}} \left(\Box \frac{\partial q}{\partial t} + \nabla \bullet F(q), \phi \right)_{\Omega^{e}} = 0$$
$$\Box = \begin{bmatrix} \mu & 0 & 0 & 0 & 0 & 0 \\ 0 & \mu & 0 & 0 & 0 & 0 \\ 0 & 0 & \mu & 0 & 0 & 0 \\ 0 & 0 & 0 & \varepsilon & 0 & 0 \\ 0 & 0 & 0 & 0 & \varepsilon & 0 \\ 0 & 0 & 0 & 0 & \varepsilon & 0 \end{bmatrix} \quad q = \begin{bmatrix} H_{x} \\ H_{y} \\ H_{z} \\ E_{x} \\ E_{y} \\ E_{z} \end{bmatrix}} \quad F(q) = -\begin{bmatrix} 0 & -E_{z} & E_{y} & 0 & H_{z} & -H_{y} \\ E_{z} & 0 & -E_{x} & -H_{z} & 0 & H_{x} \\ -E_{y} & E_{x} & 0 & H_{y} & -H_{x} & 0 \end{bmatrix}^{T}$$
Flux

Define numerical flux F* and integrate by parts: DG weak form

$$\left(\Box \frac{\partial q^{-}}{\partial t}, \phi\right)_{\Omega^{e}} = -\nabla \bullet \mathbf{F}(q^{-}), \phi_{\Omega^{e}} + \vec{n} [[\mathbf{F}(q^{-}) - \mathbf{F}^{*}(q^{-}, q^{+})], \phi_{\partial\Omega^{e}} \right)$$

$$n\Box(F - F^{*})(q, q^{+}) = \begin{cases} -Y^{+}\vec{n} \times [E^{+} - E] - \alpha \vec{n} \times \vec{n} \times [H^{+} - H] / (Y + Y^{+}) \\ Z^{+}\vec{n} \times [H^{+} - H] - \alpha \vec{n} \times \vec{n} \times [E^{+} - E] / (Z + Z^{+}) \end{cases}$$

SEDG Method

□ Spectral element discretization: tensor product basis of 1D Lagrange interpolation polynomial based on Legendre-Gauss-Lobatto grids.

$$q_{N}^{e}(x^{e}) = \sum_{i=0}^{N} \sum_{j=0}^{N} \sum_{k=0}^{N} q_{ijk} l_{i}(\xi) l_{j}(\eta) l_{k}(\gamma)$$

Hexahedral body-fitted mesh: Gorden-Hall mapping

$$x^{e} = (x, y, z) \quad \theta = (\xi, \eta, \gamma) \in [-1, 1]^{3}$$

□ 1D mass matrix is diagonal:

$$\widehat{M}_{ij} = \sum_{k=0}^{N} w_k l_i(\xi_k) \, l_j(\xi_k)$$

Surface integration:

$$\oint_{\partial\Omega} \mathbf{n} \cdot (\mathbf{F}_{\alpha} - \mathbf{F}_{\alpha}^{*}) l_{i}(\xi) l_{i}(\eta) J d\xi d\eta$$



Spectral Element DG Method

Semi-Discrete Formulations

$$M^{\mu} \frac{dH_x}{dt} + (D_y E_z - D_z E_y) = R(F_H)_x$$
$$M^{\mu} \frac{dH_y}{dt} + (D_z E_x - D_x E_z) = R(F_H)_y$$
$$M^{\mu} \frac{dH_z}{dt} + (D_x E_y - D_y E_x) = R(F_H)_z$$
$$M^{\varepsilon} \frac{dE_x}{dt} - (D_y H_z - D_z H_y) = R(F_E)_x$$
$$M^{\varepsilon} \frac{dE_y}{dt} - (D_z H_x - D_x H_z) = R(F_E)_y$$

Mass and Stiffness matrices

$$\begin{aligned}
M^{\mu} &= \mu(\psi_{ijk}, \psi_{\hat{i}\hat{j}\hat{k}}), M^{\varepsilon} = \varepsilon(\psi_{ijk}, \psi_{\hat{i}\hat{j}\hat{k}}) \\
D_{x} &= \left(\frac{\partial \psi_{ijk}}{\partial x}, \psi_{\hat{i}\hat{j}\hat{k}}\right), D_{y} = \left(\frac{\partial \psi_{ijk}}{\partial y}, \psi_{\hat{i}\hat{j}\hat{k}}\right), D_{z} = \left(\frac{\partial \psi_{ijk}}{\partial z}, \psi_{\hat{i}\hat{j}\hat{k}}\right)
\end{aligned}$$

$$(\psi_{ijk},\psi_{\hat{i}\hat{j}\hat{k}}) = J(\hat{M}\otimes\hat{M}\otimes\hat{M})$$

$$\begin{pmatrix} \frac{\partial \psi_{ijk}}{\partial x}, \psi_{\hat{i}\hat{j}\hat{k}} \end{pmatrix} = J(G^{\xi x}D_{\xi} + G^{\eta x}D_{\eta} + G^{\gamma x}D_{\gamma}),$$

$$\begin{pmatrix} \frac{\partial \psi_{ijk}}{\partial y}, \psi_{\hat{i}\hat{j}\hat{k}} \end{pmatrix} = J(G^{\xi y}D_{\xi} + G^{\eta y}D_{\eta} + G^{\gamma y}D_{\gamma}),$$

$$\begin{pmatrix} \frac{\partial \psi_{ijk}}{\partial z}, \psi_{\hat{i}\hat{j}\hat{k}} \end{pmatrix} = J(G^{\xi z}D_{\xi} + G^{\eta z}D_{\eta} + G^{\gamma z}D_{\gamma}),$$

$$\begin{split} D_{\xi} &= \hat{M} \otimes \hat{M} \otimes \hat{M} \hat{D}, \\ D_{\eta} &= \hat{M} \otimes \hat{M} \hat{D} \otimes \hat{M}, \\ D_{\gamma} &= \hat{M} \hat{D} \otimes \hat{M} \otimes \hat{M} \end{split}$$

Tensor product of 1D differentiation matrix: *matrix-matrix products*

Communication

Spectral Element DG Method

Interface

Interface

Interface

Decompose face values into
 local and neighboring element
 and assign the same form:
 (eg, central flux)

$$n_x = -n_x^+ \to n_x \frac{(E_x^+ - E_x)}{2} = \frac{-n_x^+ E_x^+ - n_x E_x}{2}$$

Communications (upwind)

$$R(F_{H}) = \begin{cases} -fH_{x} + \alpha(n_{y}E_{y} - n_{y}E_{z}) & fH_{x} = n_{z}E_{y} - n_{y}E_{z} \\ -fH_{y} + \alpha(n_{z}E_{x} - n_{x}E_{z}) & fH_{y} = n_{z}E_{x} - n_{x}E_{z} \\ -fH_{z} + \alpha(n_{x}E_{y} - n_{y}E_{x}) & fH_{z} = n_{x}E_{y} - n_{y}E_{x} \end{cases}$$

$$R(F_{E}) = \begin{cases} fE_{x} + \alpha(n_{y}H_{y} - n_{y}H_{z}) & fE_{x} = -(n_{x}H_{y} - n_{y}H_{z}) \\ fE_{y} + \alpha(n_{z}H_{x} - n_{x}H_{z}) & fE_{y} = -(n_{z}H_{x} - n_{x}H_{z}) \\ fE_{z} + \alpha(n_{x}H_{y} - n_{y}H_{x}) & fE_{z} = -(n_{x}H_{y} - n_{y}H_{z}) \end{cases}$$



Messages between element faces for Ex,Ey,Ez,Hx,Hy,Hz stored into a single array: <u>communication latency reduction by 6x</u>

Spatial Operators



Semidiscrete scheme::

$$\frac{dq_N}{dt} = -Aq_N,$$

$$q_N = (H_x^N, H_y^N, H_z^N, E_x^N, E_y^N, E_z^N)$$

Exponential time integrator (Krylov subspace approximation)

$$\begin{aligned} q_{N} (n\Delta t + \Delta t) &= e^{-\Delta tA} q_{N} (n\Delta t) \\ e^{\Delta tA} q &= q + \frac{\Delta tA}{1!} q + \frac{(\Delta tA)^{2}}{2!} q + \dots \approx K_{m} (A,q) = \operatorname{span} \{q, Aq, A^{2}q, \dots, A^{m-1}q\} \\ e^{\Delta tA} q &\approx V_{m} e^{\Delta tH} V_{m}^{T} q = V_{m} X e^{\Delta t\Lambda} X^{-1} V_{m}^{T}, e_{n+1} = \left\| e^{-\Delta tA} q_{N}^{n} - V_{m} e^{-\Delta tH_{m}} V_{m}^{T} q_{N}^{n} \right\| \leq c_{e} \Delta t^{m} \\ \left\| V_{m} e^{-\Delta tH_{m}} V_{m}^{T} \right\|_{2} \leq \left\| e^{-\Delta tH_{m}} \right\|_{2} \leq e^{\mu (-\Delta tH_{m})} \leq e^{\mu (-\Delta tA)} \leq 1 \end{aligned}$$

For any real matrix A (n x n) and the Hessenberg matrix H (m x m) is unconditionally stable if the eigenvalues of (-A) are in the negative half-plane

3D Cylindrical Waveguide Example



□ Analytic solutions:

$$H_{x} = -k_{y}w\pi\gamma^{-2}\sin(k_{x}\pi x)\cos(k_{y}\pi y)\sin(\pi t - k_{z}z)$$
$$H_{y} = k_{x}w\pi\gamma^{-2}\cos(k_{x}\pi x)\sin(k_{y}\pi y)\sin(\pi t - k_{z}z)$$
$$H_{z} = 0$$

$$E_{x} = k_{x}k_{z}\pi\gamma^{-2}\cos(k_{x}\pi x)\sin(k_{y}\pi y)\sin(\pi t - k_{z}z)$$
$$E_{y} = k_{y}k_{z}\pi\gamma^{-2}\sin(k_{x}\pi x)\cos(k_{y}\pi y)\sin(\pi t - k_{z}z)$$
$$E_{z} = \sin(k_{x}\pi x)\sin(k_{y}\pi y)\cos(\pi t - k_{z}z)$$



Exponential Integrator: Convergence in Time

Min & Fischer



Example 1/3: Nanophotonic Devices

Nanophotonics Simulations

Collaboration with CNM/ANL --> co-design proposal w/ MSD, CNM, MCS





Strong surface enhanced electric field (10-20 times larger) for nanoscale devices Energy of the electric field |E|: Intensity is proportional to $|E|^{4} > 10^{8}$

Standard FDTD does not capture correct profile on the surface (with strong oscillations): requires very dense grid resolution in order to represent high gradient fields

Electron Movement in Metal

- Harmonic oscillators
- Drude model: free electrons
- Lorentzian model: bound-electrons

Air (n=1) Au

Transmission calculations:

0.6

0.5

Transmission 8.0 8.0

0.2

0.1

400

600



Nanophotonics Simulations

Exploring a Variety of Thin Film Metal Nanoholes Nan



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Validation for Thin Film Metallic Nano Holes

Transmission Calculations: comparisons for Finite-Difference Time-Domain (FDTD) method and SEDG method.

Body-fitted SEDG mesh requires **48x** fewer gridpoints.

3D Nanohole

J. Phys: 2010, Min, Montgomery, Fischer, Gray





Accelerator Modeling



Beam Propagation Behaviors

Ultra-relativistic electron beam
 Free Space and Conducting Pipe:





Example (2/3): Moving Electron Beam

□ Collaboration with APS/ANL → ComPASS project : SLAC, Fermi, TechX.

3D moving beam in a pillbox cavity: wakefield simulations by SEDG+RK4 (PAC 2007, Min, Fischer, Chae)



Typical dimension: 10cm x 10 cm x 1m Beam size < 0.03 mm

There is a big scale difference between beam length and the whole device

Accelerator: Code Validation

Beam dynamics Simulations: Wakefield and Wake Potential Calculations

SRF2007, Min, Fischer, & Chae



Spectral-element Mesh

TESLA Cavity (Moving Window)

Accelerator: Beam Simulations

PAC 2009, Min and Fischer

Speedup with Moving Window algorithms

– **20x** reduction in cost over full-domain case



Boltzmann Equations

Boltzmann equations:

$$\frac{\partial f_{\alpha}}{\partial t} + \mathbf{e}_{\alpha} \cdot \nabla f_{\alpha} = C$$

Evolution of a single-particle distribution function f at location x with velocity at time t. Collision term C describes these effect of binary collisions on a single particle distribution with very complicated nonlinear integral form.

Discrete Boltzmann equations with BGK approximation:

simplified linearized form for the collision term

Lattice Boltzmann Method

□ Lattice Boltzmann equations: restrict the particles to have 9-velocity (2D) and 19-velocity (3D) fields: $\frac{\partial f_{\alpha}}{\partial t} + \mathbf{e}_{\alpha} \cdot \nabla f_{\alpha} = -\frac{1}{2} f_{\alpha} - f_{\alpha}^{eq}$

Lattice Boltzmann Equation (T. Lee, C-L Lin, 2001): $f_{\alpha} \mathbf{x}, t = f_{\alpha} \mathbf{x} - \mathbf{e}_{\alpha} \Delta t, t - \Delta t - \frac{1}{\lambda} \int_{t-\Delta t}^{t} f_{\alpha} - f_{\alpha}^{eq} dt'$ $f_{\alpha} \mathbf{x}, t = f_{\alpha} \mathbf{x} - e_{\alpha} \Delta t, t - \Delta t - \frac{\Delta t}{2\lambda} f_{\alpha} - f_{\alpha}^{eq} \sum_{\mathbf{x} - e_{\alpha} \Delta t, t-\Delta t} - \frac{\Delta t}{2\lambda} f_{\alpha} - f_{\alpha}^{eq} \sum_{\mathbf{x}, t-\Delta t} f_{\alpha} - f_{\alpha} - f_{\alpha$

Modified particle distribution function:

$$\overline{f_{\alpha}} = f_{\alpha} + \frac{f_{\alpha} - f_{\alpha}^{eq}}{2\tau}, \quad \overline{f_{\alpha}^{eq}} = f_{\alpha}^{eq}; \quad \rho = \sum_{\alpha=0}^{8} \overline{f_{\alpha}^{eq}} = \sum_{\alpha=0}^{8} \overline{f_{\alpha}}, \quad \rho u = \sum_{\alpha=0}^{8} e_{\alpha} \overline{f_{\alpha}^{eq}} = \sum_{\alpha=0}^{8} e_{\alpha} \overline{f_{\alpha}}$$
$$\overline{f_{\alpha}} \quad \mathbf{x}, t \quad -\overline{f_{\alpha}} \quad \mathbf{x} - \mathbf{e}_{\alpha} \Delta t, t - \Delta t \quad = -\frac{1}{\tau + 0.5} \quad \overline{f_{\alpha}} - \overline{f_{\alpha}^{eq}} \quad (\mathbf{x} - \mathbf{e}_{\alpha} \Delta t, t - \Delta t)$$
$$q_{\alpha} \quad \mathbf{x}, t \quad = \overline{f_{\alpha}} \quad \mathbf{x}, t \quad -\frac{1}{\tau + 0.5} \quad \overline{f_{\alpha}} - \overline{f_{\alpha}^{eq}} \quad (\mathbf{x}, t)$$

Collision step: $q_{\alpha} \mathbf{x}, t = \overline{f_{\alpha}} \mathbf{x}, t - \frac{1}{\tau + 0.5} \overline{f_{\alpha}} - \overline{f_{\alpha}^{eq}} (\mathbf{x}, t)$

Advection step: $\frac{\partial q_{\alpha}}{\partial t} + e_{\alpha} \Box \nabla q_{\alpha} = 0$

Min & Lee



Convergence for a Vortex Roll-Up Problem

SEDG lattice Boltzmann Method



Re=10000; E=12 x 12, N=7, 9, 11



Spectral Element DG-LBM Method

SEDG lattice Boltzmann Method

Flows Past a Cylinder

JCP 11, Min & Lee

Code Validation: Drag Coefficient

Streamline for Re= 9500



E=3,758, N=5, Total Grid Points =135,288.

Re = 9500 & Ma = 0.05



Vortex Method : P. Koumoutsakos, J. Fluid Mech., 1995 Spectral-Element Method (SEM): P. Fischer, J. Comp. Phys., 1997

Expansion in Applications and Collaborations:

- Accelerator modeling: possible participation for SciDAC ComPASS project
- Nanomaterial applications: co-design proposal on design/optimization
- Lattice Boltzmann modeling for compressible flows: with CUNY
- **Expansion to Nano Solar Cell Applications**
 - time-dependent Schrödinger-Poisson solver
- **Enhance Geometric Flexibility: hybrid element meshes**
- Efficient Timestepping Methods (e.g., SSP, exponential integrators)
- Alternative Programming Models for Extreme Scale

Thank You!